The frequency dependence of coercivity in ferroelectrics

Yifan Yuan Group meeting 20200612 Cu/P(VDF-TrFE)/Cu capacitor-type device and scaling behaviors of coercivity E_c over frequency f.



- two different regimes of switching dynamics
- domain switching in a ferroelectric involves domain wall motion
- Equivalent to the dependence upon pulse width

 $E_{\rm c} = (E_{\rm c+} - E_{\rm c-})/2$

The Ishibashi–Orihara power law $E_c \sim f^{\beta}$

KAI model for polarization switching kinetics

 $P(t) = 2P_s \left\{ 1 - \exp\left[-\left(\frac{t}{t_0}\right)^n\right] \right\}$

Under constant electric field E

The volume fraction of reversed domains, q(t) is given by (generalized formula)

$$q(t) = 1 - \exp\left[-\int_{0}^{t} C(t, t') n(t') \, \mathrm{d}t'\right], \quad (1)$$

where C(t,t') is the volume of the domain at time t which has grown from the nucleus formed at time t', and n(t') is the number of nuclei per unit volume and unit time nucleated at time t'

$$C(t,t') = C_d \left[\int_{t'}^t v(t'') \, \mathrm{d} t'' \right]^d, \qquad (2)$$

where v(t) is the velocity of domain wall, d is the growth dimension and C_d is a constant

The nucleation and growth process actually occur under an oscillating field.

$$n(t) dt = n_E(E) dE, \qquad (3)$$

The mutual overlap of growing domains

 $a(\omega t) = 1 - \exp\left(-a_{ex}(\omega t)\right), \qquad (3)$



The Ishibashi–Orihara power law $E_c \sim f^{\beta}$

(1)+(2)+(3) => (4) $q(E) = 1 - \exp\left[-\int_{0}^{E} C_{d}\left(\int_{E'}^{E} v(E'') \frac{\mathrm{d}t}{\mathrm{d}E}\Big|_{E=E''} \mathrm{d}E''\right)^{d} n_{E}(E') \mathrm{d}E'\right], \qquad (4)$

Let us apply an oscillating field E=h(ft) with a frequency f, where h(x) is a periodic function with a period of unity

$$q(E) = 1 - \exp\left[-f^{-d}\Phi(E)\right]$$
(5a)

with
$$\Phi(E) = \int_0^E C_d \left[\int_{E'}^E v(E'') \{h'(h^{-1}(E''))\}^{-1} dE'' \right]^d n_E(E') dE'$$
 (5b)

From the experimental data of TGS, $(NH_2CH_2COOH)_3 \cdot H_2SO_4$.

$$\Phi \propto E^{\alpha}$$
(6)

with α being about 5.7 both for the sinusoidal and triangular waves.

If we define the coercive field *Ec* as q(l)

The Ishibashi–Orihara power law $E_c \sim f^{\beta}$

For TGS, *d* is found, from the switching transient for the step field, to be about 1.3 - 1.6. Thus, the exponent of the coercive field is found to be about 0.23.



Why is there a cross-over coercive field?

Theoretical work on the ac dynamics of domain walls



Three regimes for domain wall motion:

- pinned (v=0)
- creep (v has an exponential relationship)
- flow (v \propto E)

Dynamic crossover between the relaxation and creep regimes (E_1)

$$\tau \approx \tau_{0} \exp(U_{B,max}/k_{B}T), \qquad (1)$$

$$U_{B,max} \approx U(E_{C0}/E)^{\mu}(1 - E/E_{C0})^{\eta}, \qquad (2)$$

$$\tau f = 1$$

$$E_{depin}/E_{C0} = [(U/k_{B}T\Lambda)(1 - E_{depin}/E_{C0})^{\eta}]^{1/\mu}, \qquad (3)$$

where $\Lambda = \ln(1/f\tau_0)$. the E_{depin} plays the role of E_1

Yang, Sang Mo, et al. *Physical Review B* 82.17 (2010): 174125.

Theoretical work on the ac dynamics of domain walls



Three regimes for domain wall motion:

- pinned (v=0)
- creep (v has an exponential relationship)
- flow ($v \propto E$)

Dynamic crossover between the creep and flow regimes (E₂)

 $v \propto \gamma E,$ $v \approx \gamma E \exp\left[-\left(U/k_B T\right) (E_{C0}/E)^{\mu} (1 - E/E_{C0})^{\eta}\right].$ (4)

$E_2/E_{C0} = [(U/k_BT)(1 - E_2/E_{C0})^{\eta}]^{1/\mu}$

Theoretical work on the ac dynamics of domain walls



The relative time duration in each regime will vary depending on f

if $f < f_{cr}$, the DW creep motion plays the major role

if $f < f_{cr}$, the viscous flow motion will dominate.

the polarization reversal (ΔP) during the first 1/4f

Theoretical work on the ac dynamics of domain walls



E1 is measured by using the E value where the switching current I started to increase significantly



 $E_2/E_{C0} = \left[(U/k_B T) (1 - E_2/E_{C0})^{\eta} \right]^{1/\mu}$

THANK YOU