

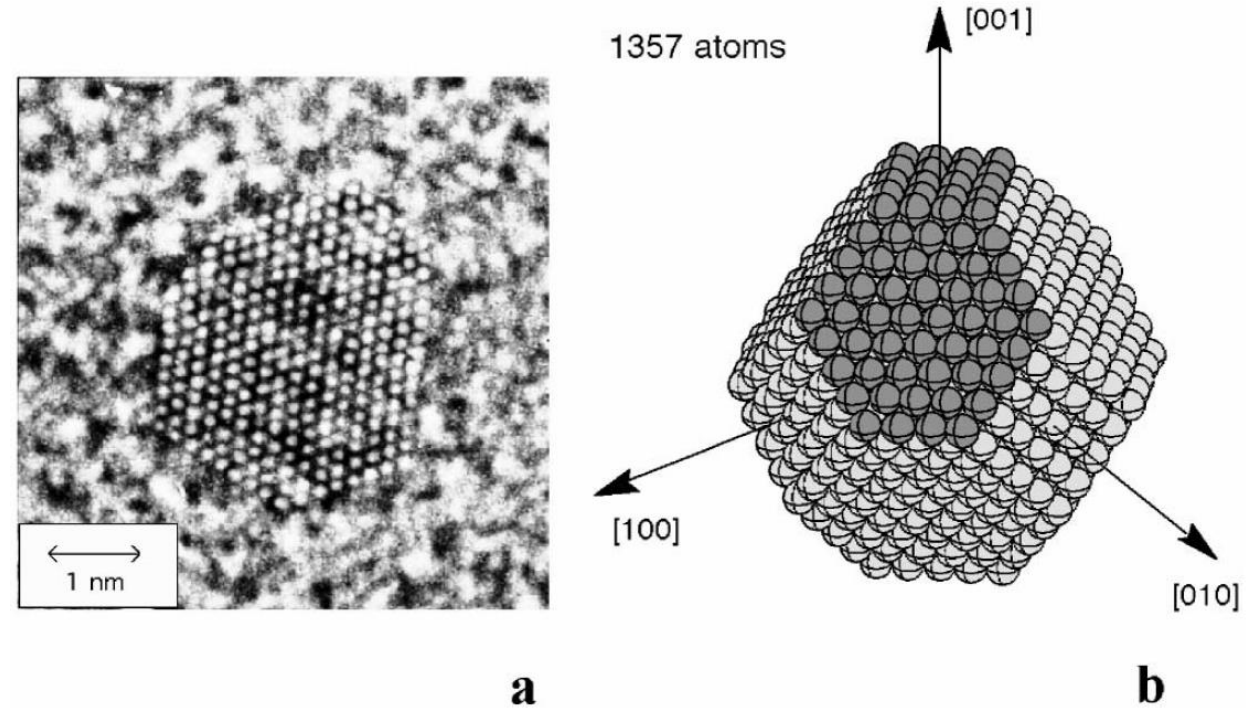
# Magnetic Anisotropy of a Single Cobalt Nanocluster

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- Haohan Wang
- Department of Physics and Astronomy
  - University of Nebraska-Lincoln

cobalt clusters (0D)

1000-3000 atoms in  
one cluster

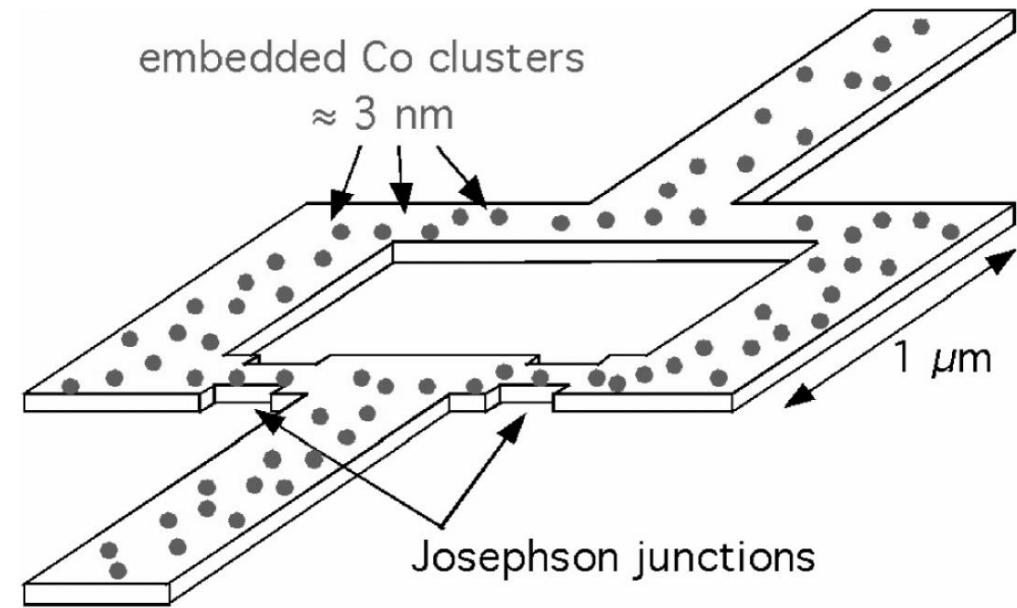


(a) High resolution transmission electron microscopy(HRTEM) observation along a  $[110]$  direction of a typical 3 nm cobalt cluster exhibiting a fcc structure.

(b) A characteristic cluster simulated for our magnetic calculations with light atoms belonging to the truncated octahedron basis and dark atoms to the  $(111)$  and  $(001)$  added facets.

Method: (LECBD) low energy cluster beam deposition

The desired sensitivity is only achieved for Co clusters embedded into the microbridges



Schematic drawing of a micro-bridge-SQUID which is patterned out of a 20-nm-thick superconducting niobium film containing a low density of 3 nm cobalt clusters (dots).

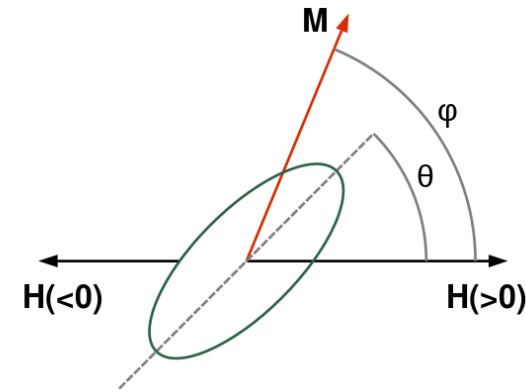
# Stoner–Wohlfarth model

The magnetization reversal is described by the potential energy

$$E(\vec{m}, \vec{H}) = E_0(\vec{m}) - \mu_0 v M_s m \cdot H \cos(\varphi)$$
$$E_0(\vec{m}) = K_u v \sin^2(\varphi - \theta)$$

where  $v$  and  $M_s$  are the magnetic volume and the saturation magnetization of the particle, respectively, and  $\mathbf{H}$  is the external magnetic field.  $E_0(\vec{m})$  is the magnetic anisotropy.  $K_u$  is anisotropy parameter.

The **switching field** is the point where the solution switches from an energy minimum ( $\partial^2 E / \partial \varphi^2 > 0$ ) to an energy maximum ( $\partial^2 E / \partial \varphi^2 < 0$ )



Dash line is easy axis

First, the magnetization of the particles is saturated in a given direction (at  $T$  35 mK).

Second, a field is applied at a temperature between 35 mK and 30 K which may or may not cause a magnetization switching.

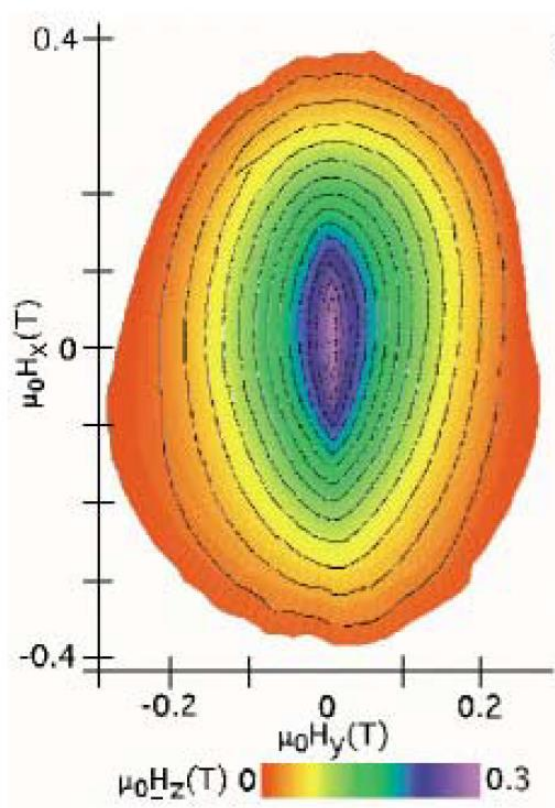
Finally, the SQUID is switched on (at  $T$  35 mK) and a field is applied in the SQUID plane to probe the resulting magnetization state.

This method can scan the entire field space.

$$E_0(\mathbf{m})/v = -K_1 m_z^2 + K_2 m_y^2 - K_4 (m_{x'}^2 m_{y'}^2 + m_{x'}^2 m_{z'}^2 + m_{y'}^2 m_{z'}^2)$$

$K_1$  and  $K_2$  are the anisotropy constants along  $z$  (easy axis) and  $y$  (hard axis).

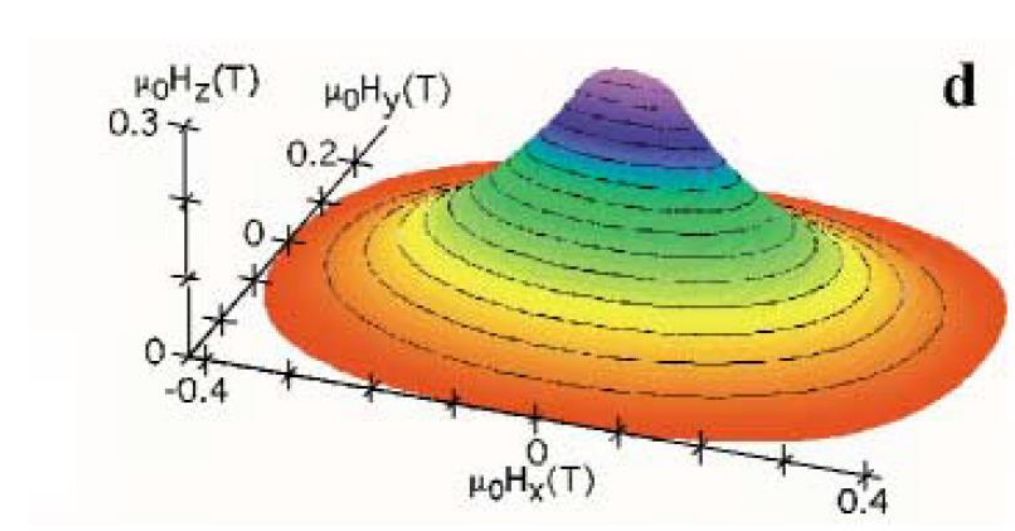
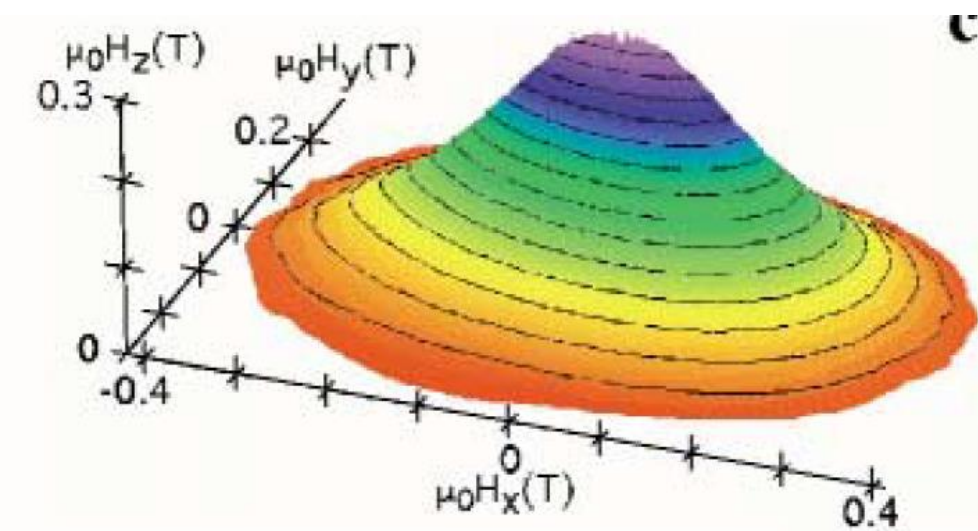
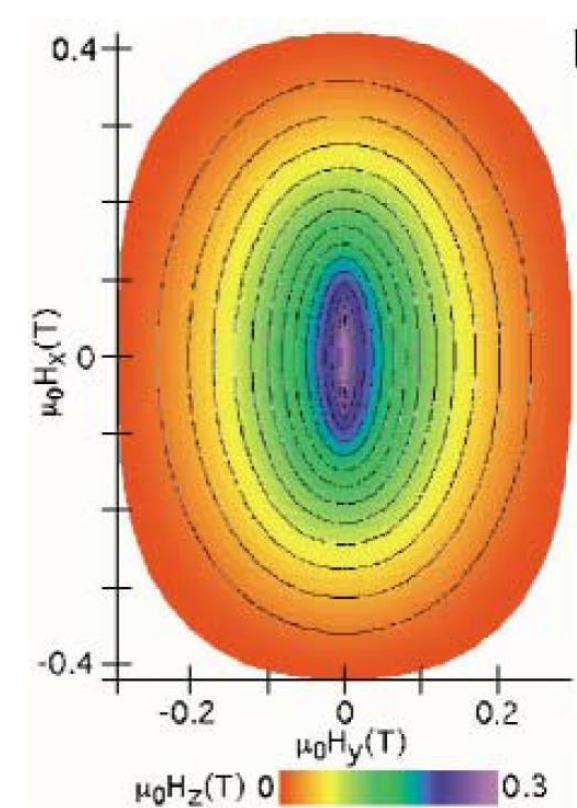
$K_4$  is the fourth order anisotropy constant, and the  $x'y'z'$  coordinate system is deduced from  $xyz$  by a 45 degree rotation around the  $z$  axis with  $z'=z$ .

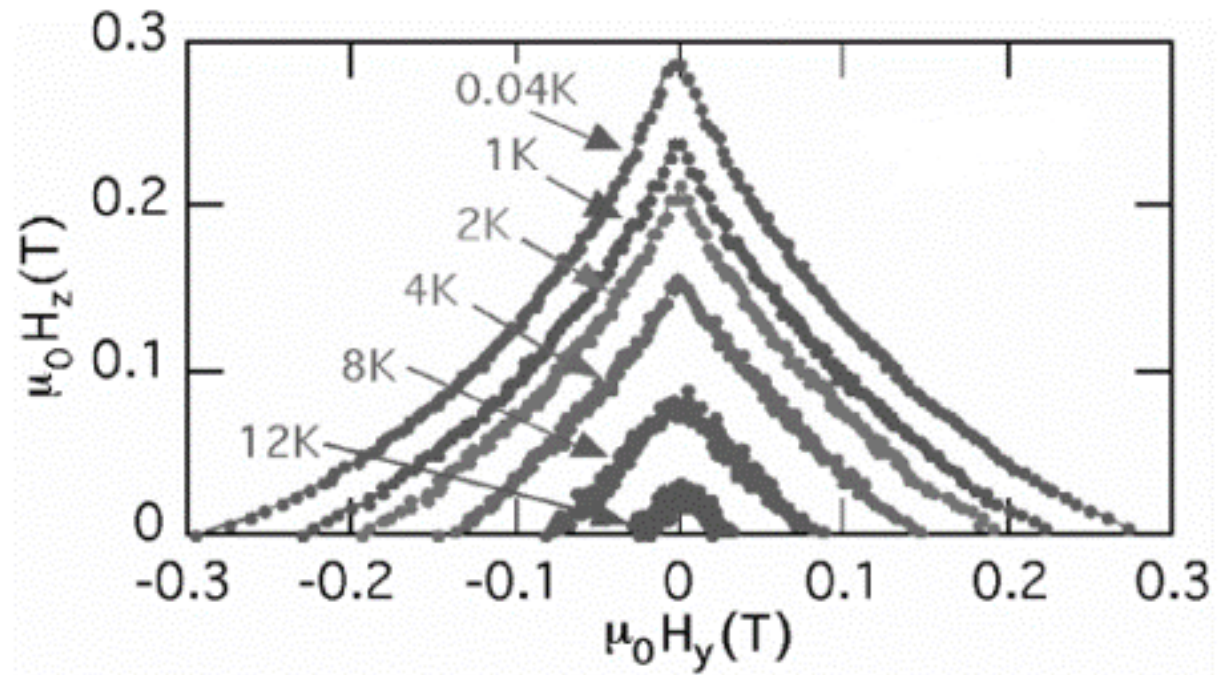


Top and side view of experiment

Top and side view of theory

	theory	experiment
$K_1(10^5\text{J/cm}^3)$	2.2	$2.0 \pm 0.3$
$K_2(10^5\text{J/cm}^3)$	0.9	$0.8 \pm 0.3$
$K_4(10^5\text{J/cm}^3)$	0.1	$0.1 \pm 0.05$





Temperature dependence of the switching field measured in the  $H_y$  -  $H_z$  plane



$$E_0(\mathbf{m}) = E_{\text{shape}}(\mathbf{m}) + E_{\text{surface}}(\mathbf{m}) \\ + E_{\text{ME}}(\mathbf{m}) + E_{\text{MC}}(\mathbf{m}).$$

$E_{\text{shape}}$  is the magnetostatic energy related to the cluster shape.

$E_{\text{surface}}$  is due to the symmetry breaking and surface strains. (main contribution)

$E_{\text{ME}}$  is the magnetoelastic anisotropy which is induced by the relaxation inside the particle if there is external stress.

$E_{\text{MC}}$  is the cubic magnetocrystalline anisotropy arising from the coupling of the magnetization with the fcc crystalline lattice as in the bulk.



# conclusion

1. Micro-SQUID technique combined with the LECBD is a powerful method to investigate the magnetic properties of nanosized magnetic particles.
2. They measure in three dimensions the switching field of individual grains, giving access to its magnetic anisotropy energy.
3. The temperature dependence of the switching field is measurable.
4. Cluster-matrix interface provides the main contribution to the magnetic anisotropy.