# Stoner Wohlfarth model in 3 dimension

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### Levels of details for ferromagnets

- Atomic level:
  - Exchange interaction that aligns atomic moments  $J_{ij}\vec{S_i}\cdot\vec{S_j}$
- Micromagnetic level
  - Smear the individual atoms into continuum, see magnetization as a function of position (domain wall)
- Domain level
  - Domains are separated by walls of zero thickness
- Nonlinear level
  - Average magnetization of the entire magnet



Saturation

#### Magnetic anisotropy

- Magneto-crystalline anisotropy
  - Microscopic
  - Single-ion
  - Symmetry of the atomic local environment
- Shape anisotropy
  - Macrocopic
  - Shape of the magnet
  - Depolarization field





#### Magnetic shape anisotropy









Attraction, low energy Polarize each other, low energy **De**polarize each other, high energy

#### Spheroidal model for anisotropy

**Mathematically**, both magneto-crystalline anisotropy and magnetic shape anisotropy can be described using the anisotropy tensor (symmetric matrix):

$$\overrightarrow{\boldsymbol{D}} = \begin{bmatrix} D_{xx} & D_{xy} & D_{xz} \\ D_{xy} & D_{yy} & D_{yz} \\ D_{xz} & D_{yz} & D_{zz} \end{bmatrix}$$

The combined matrix can be diagonalized and along the principle axis, the tensor looks like

$$\overleftarrow{\boldsymbol{D}} = \begin{bmatrix} D_{11} & 0 & 0 \\ 0 & D_{22} & 0 \\ 0 & 0 & D_{33} \end{bmatrix}.$$

Geometrically, these matrices can be described using ellipsoids.

A simplified case assumes the ellipsoid is spherioid (uniaxial).

$$D_{11} = D_{22}$$

Anisotropy energy:

$$E_A = -\vec{M} \cdot \vec{D} \cdot \vec{M} = -M^2 \sin^2 \theta D_{11} - M^2 \cos^2 \theta D_{22} = -M^2 \sin^2 \theta (D_{11} - D_{22}) - M^2 D_{22}$$

$$E_A = K \sin^2 \theta$$
,  $K = -M^2 (D_{11} - D_{22})$ 

$$\vec{M} = (\sin \theta, \cos \theta, 0)$$

## Stoner Wohlfarth model: single domain, homogeneous magnetization



Anisotropy energy:

$$E_A = K \sin^2 \theta$$
,  $K = -M^2 (D_{11} - D_{22})$ 

Zeeman energy:

 $E_Z = -HMcos(\phi - \theta)$ 

Total energy:

 $E = K\sin^2(\theta) + HM\cos(\phi - \theta)$ 

- The direction of the **magnetization** is a result of competition between the anisotropy energy and the Zeeman energy.
- Since the magnetic anisotropy is fixed for a sample, we will study the dependence of  $\vec{M}$  on H and  $\phi$ .

#### 3 D model

Assume that the angle between  $\vec{M}$  and  $\hat{A}$  is  $\alpha$ .

Assuming uniaxial, the principle axis is represented by vector  $\hat{A}(\theta_A, \phi_A)$ .



Anisotropy energy:  $E_A = K \sin^2 \alpha$  $\cos \alpha$  $= \sin \theta_A \cos \phi_A \sin \theta \cos \phi$  $+\sin\theta_A\sin\phi_A\sin\theta$ Zeeman energy:  $E_{z} = -\sin\theta\cos\phi HM$ Total energy:  $E_{total} = K(1 - \cos^2 \alpha) - \sin \theta \cos \phi HM$  $= K[(1 - \cos^2 \alpha) - \frac{HM}{K} \sin \theta \cos \phi]$ Dimensionless parameter is enough

to describe the field strength.

#### Polar map of energy $\theta = 30, \phi = 30, H = 0$ .



#### North hemisphere

South hemisphere

 $\theta = 30, \phi = 30, H = 0.$ 















#### Field dependence (MH loop)

 $\theta = 30, \phi = 30$ 



#### MH loop at different $\phi$ , $heta=10~^\circ$



#### MH loop at different $\phi$ , $\theta = 45$ °



















#### Compare to experiment



















#### Conclusion

- 3 D Stoner Wohlfarth model can be simulated at different uniaxial principle axis direction
- The in-plane and out-of-plane MH loop show different behavior.
- The out-of-plane is similar to what we have observed.